

Tests of General Relativity and Black hole spectroscopy with quasinormal modes

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1 Introduction

- The first detection - GW150914
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- Motivation
- The aim of the project

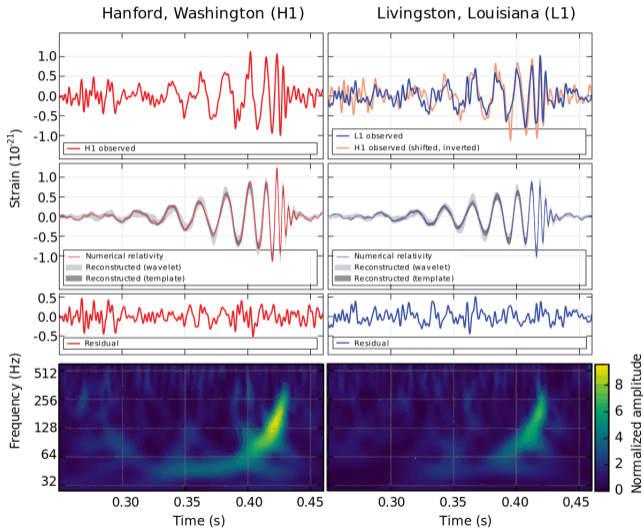
2 Scalar Perturbation

- Scalar Perturbation-Introduction
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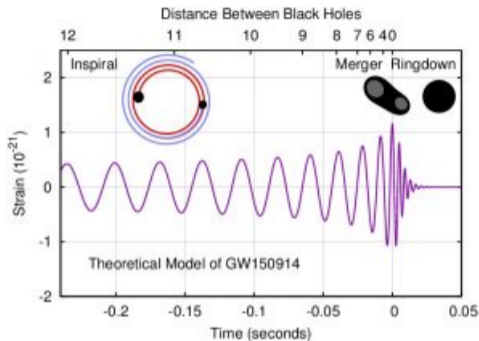
3 Methodology and Application

- WKB Method
- Application

The first detection - GW150914



Inspiral-Merger-Ringdown



- Ringdown phase is well described by perturbing the background black hole spacetime by any test field.
- It is assumed that this model is valid only when black hole is very close to reaching its equilibrium.

- **Stability of black holes:** Given the initial perturbation, will the perturbation remain bounded at all times?

[Regge, Tullio and Wheeler, John A. (*Phys. Rev* 1957)]

- **Gravitational wave Astronomy:** To extract information of the black hole like their mass, charge, and spin.
- **Tests of General Relativity:** Different modes should give us a stringent test for GR. Using different techniques within GR independently to test GR.
- **Confirmation of BHs:** Detection of these frequencies also confirm the existence of the BHs.

The aim of the project



What will happen to the black hole when it is perturbed by an external field?

- A perturbed black hole will go through quasinormal modes ringdown. The aim of the project was to calculate these quasinormal modes.
- As a toy model I will treat scalar fields as an external test field.

Scalar Perturbation-Introduction

- Klein Gordan Equation for massless scalar field in spherically symmetric and static background

$$\square\phi = 0$$

- Decomposing the function as

$$\phi(t, r, \theta, \phi) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \frac{\psi_{\ell m}^{s=0}(r)}{r} P_{\ell m}(\theta) e^{-i\omega t} e^{im\phi}$$

- By using tortoise coordinate r_* defined as $\frac{dr}{dr_*} \equiv (fh)^{1/2}$, we get the Schrodinger equation:

$$\frac{d^2\psi_l^{s=0}}{dr_*^2} + [\omega^2 - V_0] \psi_l^{s=0} = 0$$

where effective potential reads as $V_0 \equiv f \frac{\ell(\ell+1)}{r^2} + \frac{(fh)'}{2r}$

- Effective Potential $V_0^{\text{Schw}} = \left(1 - \frac{2M}{r}\right) \left(\frac{\ell(\ell+1)}{r^2} + \frac{2M}{r^3}\right)$
- **Horizon:** Classically, nothing should leave horizon, therefore

$$\phi \sim e^{-i\omega(t+r_+)}$$

- **Infinity:** Discarding unphysical waves entering from infinity.

$$\phi \sim e^{-i\omega(t-r_+)}$$

- Now our aim is to solve for ω in the radial equation satisfying the above boundary conditions and show that the eigen functions are damped in nature.

- For any massless spin ($s=0$ scalar, $s=1$ Electromagnetic, $s=2$ Gravitational):

$$\frac{d^2 \Psi_\ell^s}{dr_*^2} + Q(r_*) \Psi_\ell^s = 0$$

where $Q(r_*) = [\omega^2 - V_s]$, $V_s = f \left(\frac{\Lambda}{r^2} + \frac{2\beta}{r^3} \right)$, $\beta = 1 - s^2$, $\lambda = \ell(\ell + 1)$

- **Validity:** Due to boundary conditions, the transmitted and reflected waves should have comparable amplitudes. Therefore, turning points degenerate or very close to each other.
- In middle region, expanding Q about maxima and using boundary conditions, we get

$$n + \frac{1}{2} = -\frac{iQ_0}{(2Q_0'')^{1/2}}$$

- Using the above result to get

$$\omega_n^2 = i \left(n + \frac{1}{2} \right) \sqrt{2Q_0''} + \left(1 - \frac{2M}{r_0} \right) \left(\frac{\Lambda}{r_0^2} + \frac{2\beta M}{r_0^3} \right)$$

Application

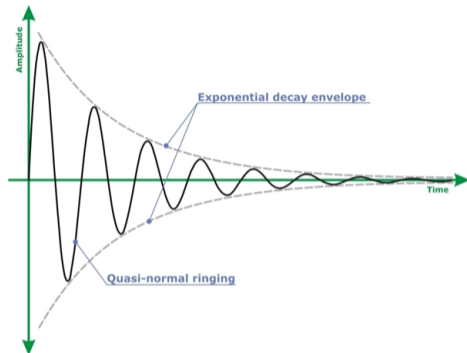
$$M = nM_{\odot}$$

For gravitational field, $s=2$

l	n	$M\omega_0 + iM\omega_j$
2	0	0.373162 - 0.0892174 <i>i</i>
	1	0.346017 - 0.274915 <i>i</i>
	2	0.302935 - 0.471064 <i>i</i>
	3	0.247462 - 0.672898 <i>i</i>
3	0	0.599265 - 0.0927284 <i>i</i>
	1	0.582355 - 0.281406 <i>i</i>
	2	0.5532 - 0.476684 <i>i</i>
	3	0.515747 - 0.677429 <i>i</i>
4	0	0.809098 - 0.0941711 <i>i</i>
	1	0.796499 - 0.284366 <i>i</i>
	2	0.773636 - 0.478974 <i>i</i>
	3	0.743312 - 0.6783 <i>i</i>

For scalar field in RN spacetime assuming $q = 0.5$

l	n	Q	$M\omega_0 + iM\omega_j$
0	0	0.5	0.247376 - 0.249018 <i>i</i>
	1	0.5	0.887109 - 0.922351 <i>i</i>
	2	0.5	1.37073 - 1.41889 <i>i</i>
	3	0.5	0.142375 - 0.863119 <i>i</i>
1	0	0.5	0.26604 - 0.145384 <i>i</i>
	1	0.5	0.700671 - 0.66721 <i>i</i>
	2	0.5	0.992813 - 0.97177 <i>i</i>
	3	0.5	0.0604634 - 0.758635 <i>i</i>



- Converting into Physical units

$$\nu = \frac{32.26}{n} (M\omega_0) \text{ kHz}, \quad \tau = \frac{n \cdot 0.4937 \cdot 10^{-5}}{(M\omega_j)} \text{ s}$$

- Bandwidth of LIGO is 10 Hz - few kHz, Corresponding mass range of BH is

$$10M_{\odot} \lesssim M \lesssim 10^3 M_{\odot}$$

- Bandwidth of LISA is $10^{-4} - 10^{-1} \text{ Hz}$, Corresponding mass range of BH is

$$1.2 \cdot 10^5 M_{\odot} \lesssim M \lesssim 1.2 \cdot 10^8 M_{\odot}$$

- Test no-hair theorem and stability of black holes
- In eikonal limit, we can relate QNM to Lyapunov exponent

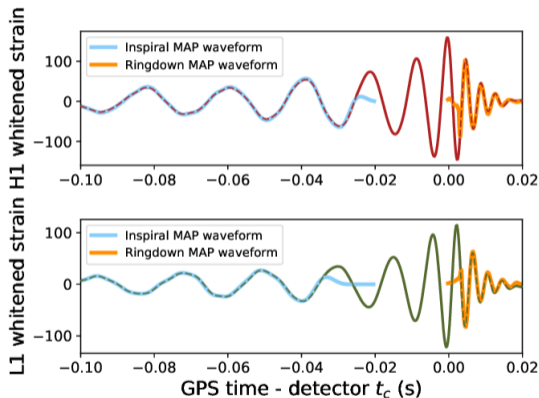
$$\omega_{QNM} = \Omega_c l + i \left(n + \frac{1}{2} \right) |\lambda|$$

where λ is the Lyapunov exponent (measures of the rate at which the trajectories diverge) of the null orbit.

Ω_c is the orbital angular velocity of the null orbit.

Open Question

- Can we quantify the beginning of the ringdown through observations?



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Thank You

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